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ABSTRACT

Ground-state rotational spectra were observed for ten symmetric-top isotopologues $H_3N\cdots S{=}C{=}S$, $H_3N\cdots{}^{34}S{=}C{=}S$, $H_3N\cdots S{=}C{=}{}^{34}S$, $H_3N\cdots S{=}{}^{13}C{=}S$, $H_3{}^{15}N\cdots S{=}C{=}S$, $H_3{}^{15}N\cdots{}^{34}S{=}C{=}S$, $H_3{}^{15}N\cdots S{=}C{=}{}^{34}S$, $H_3{}^{15}N\cdots S{=}{}^{13}C{=}S$, $H_3{}^{15}N\cdots{}^{33}S{=}C{=}S$, and $H_3{}^{15}N\cdots S{=}C{=}{}^{33}S$, the first five in their natural abundance in a mixture of ammonia and carbon disulphide in argon and the second group with enriched $^{15}NH_3$. The four asymmetric-rotor isotopomers $H_2DN\cdots S{=}C{=}S$, $H_2DN\cdots{}^{34}S{=}C{=}S$, $H_2DN\cdots S{=}C{=}{}^{34}S$, and $HD_2N\cdots S{=}C{=}S$ were investigated by using a sample composed of ND_3 mixed with CS_2. Rotational constants, centrifugal distortion constants, and ^{33}S nuclear quadrupole coupling constants were determined from spectral analyses and were interpreted with the aid of models of the complex to determine its symmetry, geometry, one measure of the strength of the intermolecular binding, and information about the subunit dynamics. The complex has C_{3v} symmetry, with nuclei in the order $H_3N\cdots S{=}C{=}S$, thereby establishing that the non-covalent interaction is a chalcogen bond involving the non-bonding electron pair of ammonia as the nucleophile and the axial region near one of the S atoms as the electrophile. The small intermolecular stretching force constant k_σ = 3.95(5) N m^{-1} indicates a weak interaction and suggests the assumption of unperturbed component geometries on complex formation. A simple model used to account for the contribution of the subunit angular oscillations to the zero-point motion leads to the intermolecular bond length $r(N\cdots S)$ = 3.338(10) Å.

I. INTRODUCTION

A wide range of non-covalent, pairwise interactions of molecules in the gas phase has been investigated spectroscopically in the last 50 years. The most important of these is the hydrogen bond because of its ubiquity in Chemistry and Biology. The halogen bond has become of increasing interest since the 1990s, and its properties have been shown[1,2] to parallel those of the hydrogen bond in many respects. Modern definitions[3,4] of both these important interactions have been published recently. Tetrel, pnictogen, chalcogen, and coinage-metal bonds are less well known non-covalent interactions and were given names only during the last seven years,[5–8] even though they had been characterised extensively in the gas phase during the last 50 years.[9] These newer types of pairwise interaction can be written in general as B···E–R, where E is an atom of group 14, 15, 16, or 11, respectively, R is the remainder of the molecule E–R, and B is a Lewis base. The atom E is the

most electrophilic region of E–R and interacts with a nucleophilic region of B, usually a non-bonding (n) electron pair or a π-bonding electron pair. As an example, the C atom is the most electrophilic region of CO_2 and interacts with the n-pair of NH_3 to form a T-shaped complex in which the linear CO_2 molecule acts as the bridge of T while the C_3 symmetry axis of NH_3 forms the stem.[10] Thus, the weak intermolecular bond is a tetrel bond.

Carbon disulphide (CS_2) is an isostructural congener of CO_2 and therefore might be expected also to form tetrel bonds in which its C atom again acts as the electrophilic region when interacting with Lewis bases such as NH_3. Examination of the molecular electrostatic surface potentials (MESPs) of CO_2 and CS_2 calculated at the MP2/6-311++G** level with Spartan[11] and shown in Fig. 1 reveals that the electrophilic (blue, most positive potential) region is differently placed in these two molecules, however. CO_2 has an electrophilic (blue) band around the C atom, while CS_2 has no such band, but instead its most electrophilic region lies near each S atom along the C_∞ axis. This indicates that CS_2 and NH_3 are likely to form a complex $H_3N\cdots S{=}C{=}S$ of C_{3V} symmetry. To the best of our knowledge, the only complex of CS_2 investigated by means of its rotational spectrum is that with H_2O, for which Ogata and Lovas[12] report that the most probable geometry is $H_2O\cdots S{=}C{=}S$, planar with C_{2V} symmetry.

We report here a detailed investigation, by chirped-pulse broadband microwave spectroscopy, of the ground-state rotational spectra of fourteen isotopologues of a complex formed by ammonia with carbon disulphide. Interpretation of the spectroscopic constants thereby determined in the light of a simple model of the complex, allows the symmetry, the order of the atoms, the geometry, the strength of binding (as measured by the intermolecular stretching force constant k_σ), and information about the subunit dynamics to be derived. The results can then be compared with those from good quality *ab initio* calculations, as recently published.[13]

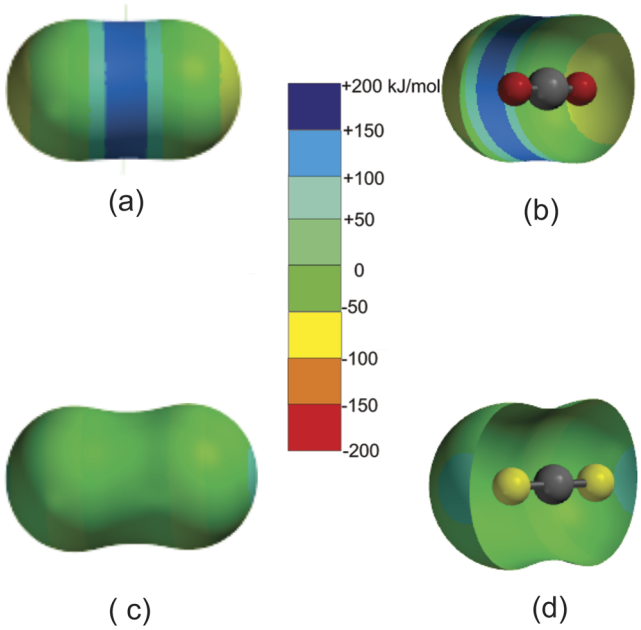

FIG. 1. Molecular electrostatic surface potentials (MESPs) on the 0.002 e/bohr^3 isosurface of carbon dioxide (a) and (b) and carbon disulphide (c) and (d) calculated at the MP2/6-311++G** level using SPARTAN. In (b) and (d), each surface has been rotated by several degrees about the C_2 axis and the front half removed.

II. EXPERIMENTAL

The chirped-pulse, Fourier-transform microwave (CP-FTMW) spectrometer used at Newcastle University has been described in detail elsewhere.[14] The spectrometer can be configured to operate in either the 2.0–8.0 or 7.0–18.5 GHz frequency bands during individual experiments. Spectra were recorded across each of these bands during the present work.

The generation of $CS_2\cdots NH_3$ was achieved by mixing CS_2 and NH_3 within an argon buffer gas and pulsing this sample through a supersonic valve (Parker, General Series 9) into a vacuum chamber at a backing pressure of 7 bars. Each of CS_2 and NH_3 were present in the gas sample at low concentrations of 1% and 2%, respectively. Isotopically enriched samples of $^{15}NH_3$ or ND_3 were used to permit experiments on isotopologues containing ^{15}N or D, respectively. Coherent polarisation of the molecular ensemble is induced by a microwave pulse (of duration 1 μs) that sweeps from high to low frequency. Propagation of this microwave pulse is perpendicular to the direction of travel of molecules introduced from the valve. The subsequent relaxation of the molecular ensemble was recorded and digitised by a 100 GS/s oscilloscope (Tektronix DPO72304SX) in the form of a free induction decay (FID) of duration 20 μs.

The sequence of polarisation pulse followed by free induction decay is complete within less than 25 μs. Molecules in the expanding jet are spatially positioned to interact with the microwave pulse for longer than 200 μs. The time scale thus allows for eight iterations of the spectral acquisition per valve pulse. An arbitrary waveform generator, phase-locked dielectric resonant oscillator, and oscilloscope are locked to a Rb-clock, which provides a 10 MHz reference that allows coherent averaging of the data in the time domain. The FIDs are co-added in the time domain prior to fast Fourier transformation (FFT) using a Kaiser-Bessel window function to generate the frequency domain spectrum. The resulting spectra have linewidths of 100 kHz (FWHM), consistent with standard deviations of approximately 10 kHz achieved when fitting transitions to the predictions of a model Hamiltonian.

III. RESULTS

A. Observed rotational spectra

The rotational spectra of the isotopologues $H_3N\cdots S{=}C{=}S$, $H_3N\cdots{}^{34}S{=}C{=}S$, $H_3N\cdots S{=}C{=}{}^{34}S$, $H_3N\cdots S{=}{}^{13}C{=}S$, $H_3{}^{15}N\cdots S{=}C{=}S$, $H_3{}^{15}N\cdots{}^{34}S{=}C{=}S$, $H_3{}^{15}N\cdots S{=}C{=}{}^{34}S$,

$H_3{}^{15}N\cdots S{=}{}^{13}C{=}S$, $H_3{}^{15}N\cdots{}^{33}S{=}C{=}S$, and $H_3{}^{15}N\cdots S{=}C{=}{}^{33}S$ (absence of a superscript mass number implies the most abundant nuclide) were observed, either in their natural abundance by using an isotopically normal mixture of ammonia and carbon disulphide or a mixture containing an isotopically enriched sample of $^{15}NH_3$. Each spectrum consisted of rotational transitions of the type $J+1 \leftarrow J$, $K \leftarrow K$ expected for the vibrational ground-state of a prolate symmetric-top molecule. Only transitions having $K = 0$ and $K = 1$ had a detectable intensity presumably because $K = 2$ states lose their room temperature population by transfer to $K = 1$ states during the supersonic expansion of the gas mixture, while transfer of population from $K = 1$ to $K = 0$ states by collision is spin-forbidden. In the ^{14}N-containing species, each $J+1 \leftarrow J$, $K \leftarrow K$ transition exhibited a hyperfine pattern characteristic of coupling of the ^{14}N nuclear spin vector $\mathbf{I}_N$ to the framework angular momentum vector $\mathbf{J}$ via the interaction of the ^{14}N nuclear electric quadrupole moment with the electric field gradient (efg) at that nucleus. Figure 2 shows a recording of the $J+1 \leftarrow J = 1 \leftarrow 0$, $K \leftarrow K = 0 \leftarrow 0$ transition of the parent isotopologue of $H_3N\cdots S{=}C{=}S$ in which the three-component hyperfine structure is resolved. No hyperfine structure was observed in transitions of the ^{15}N-containing molecules, except for those of $H_3{}^{15}N\cdots{}^{33}S{=}C{=}S$ and $H_3{}^{15}N\cdots S{=}C{=}{}^{33}S$, which both carried nuclear quadrupole hyperfine structure characteristic of molecules containing a single ^{33}S nucleus (spin quantum number $I = 3/2$).

The observed frequencies for each isotopologue were accordingly fitted in an iterative nonlinear least squares procedure in which the Hamiltonian given in Eq. (1) was constructed in the coupled symmetric rotor basis $\mathbf{J} + \mathbf{I}_{N(orS)} = \mathbf{F}$ and diagonalized in blocks of the quantum number F by using Western's program PGOPHER.[15] The form of the Hamiltonian H used was

$$H = H_R - \frac{1}{6}\mathbf{Q_X} : \nabla\mathbf{E_X}, \quad (1)$$

in which H_R is the familiar rotational energy operator for a semi-rigid prolate symmetric-top molecule. It was necessary to include only the quartic terms $-D_J\mathbf{J}^4$ and $-D_{JK}\mathbf{J}^2J_a{}^2$ in H_R for a fit commensurate with the estimated error of frequency measurement. In the second term in Eq. (1), $\mathbf{Q_X}$ is the electric quadrupole moment dyadic of nucleus X (= ^{14}N or ^{33}S) and $\nabla\mathbf{E_X}$ is the electric field gradient dyadic at nucleus X. Since the X nucleus lies on the unique axis a of the prolate symmetric-top molecule, the only independent element of the quadrupole coupling tensor is $\chi_{aa}(X) = eq_{aa}Q_X$ of the X nucleus, where $q_{aa} = \partial^2V/\partial a^2$, Q_X is the conventional electric quadrupole moment of nucleus X, and e is the proton charge. The spectroscopic constants determined in the final cycle of the least-squares fit for each of the 10 symmetric-top isotopologues are shown in Table I.

Spectroscopic constants determined for the four asymmetric-rotor isotopologues $H_2D^{14}N\cdots{}^{32}S{=}{}^{12}C{=}{}^{32}S$, $H_2D^{14}N\cdots{}^{34}S{=}{}^{12}C{=}{}^{32}S$, $H_2D^{14}N\cdots{}^{32}S{=}{}^{12}C{=}{}^{34}S$, and $HD_2{}^{14}N\cdots{}^{32}S{=}{}^{12}C{=}{}^{32}S$ are included in Table I. In each case, only $K_{-1} = 0$ transitions could be observed, a familiar consequence of the breaking of C_{3v} symmetry and thereby relaxation of the $K+1 \rightarrow K = 1 \rightarrow 0$ spin-forbidden collisional propensity rule mentioned earlier. Because the $K_{-1} = 1$ levels are several cm^{-1} higher in wavenumber than the $K_{-1} = 0$ levels in the D species, they are susceptible to cooling in the expansion. The ^{14}N nuclear quadrupole hyperfine structure was insufficiently resolved for these species, and each spectrum was fitted by using the appropriate asymmetric-rotor Hamiltonian to give the only determinable spectroscopic constants $(B+C)/2$ and Δ_J, as included in Table I. Observed frequencies and the spectroscopic constants determined in the final cycle of the PGOPHER fits for all isotopologues investigated are available in the supplementary material.

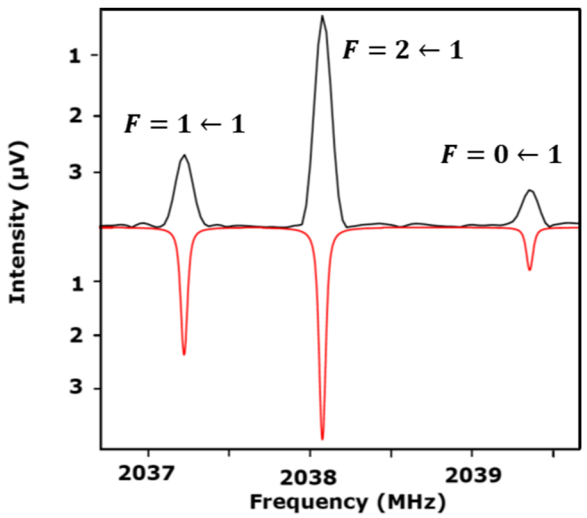

FIG. 2. The $J+1 \leftarrow J = 1 \leftarrow 0$ transition of the parent isotopologue of $H_3N\cdots S{=}C{=}S$ showing the three ^{14}N nuclear quadrupole hyperfine components. This was recorded by collecting 1.15×10^6 free induction decays. The red spectrum is that simulated with PGOPHER by using the spectroscopic constants given in Table I.

B. Experimentally determined distance r(N···S)

The magnitude of the changes in the rotational constants B_0 of the parent species on isotopic substitution at each atom establishes that the order of the atoms is $H_3{}^{14}N\cdots{}^{32}S{=}{}^{12}C{=}{}^{32}S$. The nature of the rotational spectra observed shows that the complex is a symmetric-top molecule. The fact that the interaction between the two molecules is very weak (see Sec. III D) suggests that we assume the r_0 geometries of ammonia and carbon disulfide are unchanged on complex formation when determining a r_0 value of the distance $r(N\cdots S)$ in the complex. The r_0 geometry of NH_3 was determined by fitting the ground-state rotational constants[16] B_0 and C_0 of $^{14}NH_3$ and $^{15}NH_3$ isotopologues, while that of CS_2 was similarly obtained from a fit of the ground-state rotational constants of several isotopologues measured in a recent high-resolution examination of the Raman spectrum.[17] In each case, Kisiel's program STRFIT[18] was employed and the results are in Table II. The observed ^{14}N nuclear quadrupole

TABLE I. Experimentally determined spectroscopic constants for all observed isotopologues of $H_3N\cdots S{=}C{=}S$.

Isotopologue	B_0 (MHz)	D_J (kHz)	D_{JK} (MHz)	$\chi(^{14}N)$ (MHz)	$\chi(^{33}S)$ (MHz)	N[a]	σ_{RMS}[b] (kHz)
$H_3N\cdots S{=}C{=}S$	1018.968 00(25)[c]	0.6688(25)	0.295 76(20)	−2.8459(37)	...	55	7.5
$H_3N\cdots{}^{34}S{=}C{=}S$	1017.259 11(24)	0.6603(20)	0.294 77(21)	−2.8487(55)	...	29	5.5
$H_3N\cdots S{=}{}^{13}C{=}S$	1017.342 86(49)	0.6683(37)	0.294 75(38)	−2.801 (27)	...	18	8.6
$H_3N\cdots S{=}C{=}{}^{34}S$	995.274 80(35)	0.6319(28)	0.283 26(30)	−2.8493(88)	...	27	7.6
$H_3{}^{15}N\cdots S{=}C{=}S$	988.089 09(33)	0.6390(27)	0.275 51(25)	...	...	17	5.8
$H_3{}^{15}N\cdots{}^{34}S{=}C{=}S$	986.680 68(39)	0.6331(36)	0.274 82(32)	...	...	16	6.8
$H_3{}^{15}N\cdots S{=}{}^{13}C{=}S$	986.408 21(55)	0.6556(44)	0.268 27(31)	...	...	8	4.7
$H_3{}^{15}N\cdots S{=}C{=}{}^{34}S$	965.040 52(57)	0.6082(54)	0.265 41(25)	...	...	14	7.9
$H_3{}^{15}N\cdots{}^{33}S{=}C{=}S$	987.377 38(68)	{0.6331(36)}[d]	0.276 94(50)	...	−16.135(76)	18	14.2
$H_3{}^{15}N\cdots S{=}C{=}{}^{33}S$	976.289 88(68)	{0.6082(54)}[d]	0.266 37(97)	...	−16.495(29)	22	14.1
$H_2DN\cdots S{=}C{=}S$	981.558 9(12)[e]	0.646(10)	...	...	...	6	14.1
$HD_2N\cdots S{=}C{=}S$	947.493 8(13)	0.613(11)	...	...	...	6	14.9
$H_2DN\cdots{}^{34}S{=}C{=}S$	980.806 8(19)	0.611(16)	...	...	...	5	19.1
$H_2DN\cdots S{=}C{=}{}^{34}S$	958.735 5(12)	0.604(10)	...	...	...	6	14.4

[a]Number of frequencies in the fit.
[b]Root mean square deviation of the fit.
[c]Numbers in parentheses are one standard deviation in units of the last significant figure.
[d]Value for D_J fixed at that for $H_3{}^{15}N\cdots{}^{34}S{=}C{=}S$ and $H_3{}^{15}N\cdots S{=}C{=}{}^{34}S$ isotopologue, respectively. Only three J available for $H_3{}^{15}N\cdots{}^{33}S{=}C{=}S$ and $H_3{}^{15}N\cdots S{=}C{=}{}^{33}S$, and three J dependent terms cannot be well defined.
[e]$(B + C)/2$ for asymmetric-top isotopologues containing D.

coupling constant of $^{14}NH_3$ is also included in Table II for later use.[19]

A fit of $r(N\cdots S)$ to the moments of inertia I_b^0 of the 10 symmetric-top isotopologues with STRFIT by assuming unchanged monomer geometries (as given in Table II) leads to the very precise value r_0 (N···S) = 3.3292(1) Å. Inclusion of the D-substituted species gives 3.3288(2) Å for this distance. An r_s version[20] is available through the isotopic substitutions at N and S. The resulting value r_s (N···S) = 3.311(2) Å. Such analyses do not take into account that the complex is very floppy in the zero-point state, as indicated by a small intermolecular stretching force constant of only ~4.0 N m^{-1} (see Sec. III D), which is of similar magnitude to those of inert gas complexes.

It is possible to account for the most serious effects of the zero-point motion of the complex $H_3N\cdots S{=}C{=}S$, namely, the effects of the intermolecular bending modes, when determining the distance $r(N\cdots S)$ from ground-state moments of inertia. The model of the complex[21] shown schematically in Fig. 3 allows for the contribution of these modes but not the intermolecular stretching mode. Each subunit NH_3 and S=C=S is assumed to execute the indicated angular oscillations, assumed axially symmetric (two-dimensionally isotropic), with respect to its centre of mass, but with the distance r_{cm} between the mass centres fixed (i.e., no intermolecular stretching). In Fig. 3, θ is the angle between the line r_{cm} and the C_3 axis of NH_3 and ϕ is the corresponding angle made with the C_∞ axis of CS_2. By using this model, it can be shown[21] that I_b^0 of the complex is related to the ground-state principal moments of inertia $I_b^{NH_3}$, $I_c^{NH_3}$, and $I_b^{CS_2}$ of the components (as calculated from their observed rotational constants recorded in Table II) by

TABLE II. Some molecular properties of ammonia and carbon disulphide.

Molecule	B_0 (MHz)	C_0 (MHz)	r_0 geometry	$\chi_0(^{14}N)$ (MHz)	$\chi_0(^{33}S)$ (MHz)
$^{14}NH_3$	298 115.37[a]	187 405[b]	r_0(N–H) = 1.015 57 Å[c] ∠(HNH) = 107.277°	−4.089 83(2)[d]	...
$^{15}NH_3$	297 388.12[a]	187 405[b]	...	...	
$^{32}S{=}{}^{12}C{=}{}^{32}S$	3271.5170(7)[e]	3271.5170(7)[e]	r_0(C=S) = 1.554 27(1)[f] Å	...	−17.1(6)[g]

[a]Reference 16.
[b]Calculated from the r_0 geometry.
[c]Calculated by fitting B_0 values of $^{14}NH_3$ and $^{15}NH_3$.
[d]Reference 19.
[e]Reference 17.
[f]Calculated by fitting B_0 values of six isotopologues of S=C=S given in Ref. 17.
[g]Calculated *ab initio* (see Sec. III C for details).

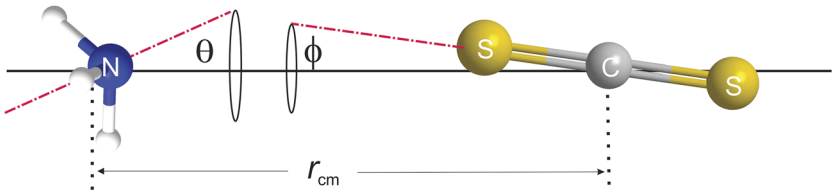

FIG. 3. Definition of r_{cm} and the oscillation angles θ and ϕ used in interpretation of the observed principal moments of inertia I_b^0 of 10 symmetric-top isotopologues of $H_3N\cdots S{=}C{=}S$ to give r_{cm} and $r(N\cdots S)$.

$$I_b^0 \approx \mu r_{cm}^2 + \frac{1}{2}I_b^{NH_3}\langle 1+\cos^2\theta\rangle + \frac{1}{2}I_c^{NH_3}\langle \sin^2\theta\rangle + \frac{1}{2}I_b^{CS_2}\langle 1+\cos^2\phi\rangle. \quad (2)$$

In Eq. (2), the angular brackets indicate averages taken over the zero-point motion and $\mu = m^{NH_3}m^{CS_2}/\left(m^{NH_3}+m^{CS_2}\right)$ is the reduced mass for the axially symmetric motion of the subunits.

It is necessary to estimate values of $\theta_{av} = \cos^{-1}\langle\cos^2\theta\rangle^{\frac{1}{2}}$ and $\phi_{av} = \cos^{-1}\langle\cos^2\phi\rangle^{\frac{1}{2}}$ for use in Eq. (2). In Sec. III C, it is shown that the angle $\theta_{av} = 27(3)^\circ$ may be determined from the ^{14}N nuclear quadrupole coupling constant $\chi(^{14}N)$, after correction for the change in efg when the CS_2 subunit is brought from infinity to the separation r_{cm}. The change in efg was estimated by calculations for NH_3 and $H_3N\cdots S{=}C{=}S$ at the B3LYP/aug-cc-pV6Z level of theory using the Gaussian Electronic Structure program.[22] A detailed analysis of calculations at the B3LYP/aug-cc-pV*n*Z (*n* = 3, 4, 5, or 6) level for free CS_2 (see Sec. III C) leads to an estimate of $\chi_0(^{33}S)$ = –17.1(6) MHz, which when combined in the way discussed in Sec. III C with the measured values of $\chi(^{33}S)$ = –16.495(46) MHz and –16.135(76) MHz for $H_3{}^{15}N\cdots S{=}C{=}{}^{33}S$ and $H_3{}^{15}N\cdots{}^{33}S{=}C{=}S$, respectively, leads to ϕ_{av} = 10(3)°. Values of r_{cm} for the ten symmetric-top isotopologues of $H_3N\cdots S{=}C{=}S$ investigated here are given in Table III, which also includes distances $r(N\cdots S)$ estimated by means of

$$r(N\cdots S) = r_{cm} - r - r', \quad (3)$$

TABLE III. Some properties of the complex $H_3N\cdots S{=}C{=}S$ derived from the rotational spectra of ten symmetric-top isotopologues.

Isotopologue	r_{cm} (Å)[a]	$r(N\cdots S)$ (Å)[a]	k_σ (N m^{-1})[b]	$10^{21}\,k_{\theta\theta}$ (J)[c]
$H_3N\cdots S{=}C{=}S$	4.959(10)	3.338(10)	3.95(2)	4.600(3)
$H_3N\cdots S{=}C{=}{}^{34}S$	4.999(10)	3.338(10)	3.91(2)	4.583(5)
$H_3N\cdots S{=}{}^{13}C{=}S$	4.959(10)	3.338(10)	3.95(2)	4.600(6)
$H_3N\cdots{}^{34}S{=}C{=}S$	4.920(10)	3.339(10)	3.95(1)	4.600(14)
$H_3{}^{15}N\cdots S{=}C{=}S$	4.954(10)	3.337(10)	4.01(2)	4.643(5)
$H_3{}^{15}N\cdots S{=}C{=}{}^{34}S$	4.994(10)	3.337(10)	3.93(3)	4.599(4)
$H_3{}^{15}N\cdots S{=}{}^{13}C{=}S$	4.954(10)	3.337(10)	3.89(3)	4.751(5)
$H_3{}^{15}N\cdots{}^{34}S{=}C{=}S$	4.915(10)	3.338(10)	4.00(2)	4.642(7)
$H_3{}^{15}N\cdots{}^{33}S{=}C{=}S$	4.934(10)	3.338(10)	...	4.613(8)
$H_3{}^{15}N\cdots S{=}C{=}{}^{33}S$	4.974(10)	3.337(10)	...	4.689(18)

[a]Calculated by using Eq. (2).
[b]Calculated by using Eq. (5).
[c]Calculated by using Eq. (6).

where *r* is the distance of N from the NH_3 mass centre and *r*′ is the distance of the inner S atom from the mass centre of CS_2, as calculated from the r_0 geometries given in Table II. The mean value of $r(N\cdots S)$ = 3.338(10) Å is longer than that obtained in the simple r_0 fit by 0.009 Å and represents the best available estimate of this quantity. As expected, it is longer than the equilibrium value r_e = 3.304 Å obtained recently[13] from an *ab initio* calculation at the CCSD(T)/aug-cc-pVTZ level [CCSD(T) = coupled cluster, doubles and perturbative triples]. The r_s value 3.311(2) Å is closer to the *ab initio* value, but the agreement might be fortuitous.

C. Interpretation of the ^{14}N and ^{33}S nuclear quadrupole coupling constants

The determination of the angles θ_{av} and ϕ_{av}, as defined in Sec. III B, relies on interpretations of the observed ^{14}N and ^{33}S nuclear quadrupole coupling constants, respectively. In each case, the observed coupling constant of the complex is first corrected for the additional efg at the nucleus X in question arising from the presence of the other subunit when in the complex. This is assumed to be a small effect and can be calculated as set out below. Then, the nuclear quadrupole coupling constant so corrected $\chi_{aa}^{corr}(X)$ is related to its free molecule value $\chi_0(X)$ by the familiar expression

$$\chi_{aa}^{corr}(X) = \frac{1}{2}\chi_0(X)\langle 3\cos^2\tau - 1\rangle \quad (\tau = \theta \text{ or } \phi), \quad (4)$$

from which $\langle\cos^2\tau\rangle$ is available. We deal with the case $\tau = \phi$ first.

Unfortunately, a value of $\chi_0(^{33}S)$ for the free CS_2 molecule is unavailable from the usual source because CS_2 has no pure rotational spectrum. Measurements of NMR spectra of liquid CS_2 have yielded the values –13.8 MHz (Ref. 23) and –14.9(3) MHz,[24] and an early *ab initio* calculated value is available.[25] Here, we employ the following approach. Density functional calculations carried out at the B3LYP/aug-cc-pV*n*Z level of theory (*n* = 3, 4, 5, and 6) with the Gaussian electronic structure package[22] were used to evaluate the ^{33}S nuclear quadrupole coupling constant $\chi_0(^{33}S)$ for each of the three molecules CS, OCS, and CS_2. For each *n*, the basis function aug-cc-pV(*n* + d) Z was used for S because S requires a tight d-function. The results for the different basis sets are recorded in Table IV. The zero-point values of this coupling constant are known to be highly accurate for CS[26] and OCS[27] and are included in Table IV. It is noted from Table IV that from *n* = 5 to *n* = 6 the calculated coupling constants change by 0.03 MHz, 0.11 MHz, and 0.17 MHz, for CS, CS_2, and OCS,

TABLE IV. Calculated values of the ^{33}S nuclear quadrupole coupling constant of CS_2, CS, and OCS at the B3LYP /aug-cc-pV(n + d)Z level of theory for n = 3, 4, 5, and 6 and its estimated value for the zero-point state of CS_2.

Molecule	Basis function aug-cc-pVnZ for B3LYP calculation	$\chi_{calc.}(^{33}S)$ (MHz)	$\chi_{obs.}(^{33}S)$ (MHz)	$\chi_{obs.}(^{33}S)/\chi_{calc.}(^{33}S)$	Est. value $\chi_{obs.}(^{33}S)$ (MHz) for CS_2 from CS ratio in col. 5	Est. value $\chi_{obs.}(^{33}S)$ (MHz) for CS_2 from OCS ratio in col. 5
CS_2	n = 6	−16.845	...	...	−17.750	−16.520
	n = 5	−16.958	...	...	−17.819	−16.534
	n = 4	−14.934	...	...	−16.471	−16.299
	n = 3	−14.966	...	...	−15.442	−16.162
CS	n = 6	+12.172	+12.8256(2)[a]	1.0537		
	n = 5	+12.206		1.0508		
	n = 4	+11.629		1.1029		
	n = 3	+12.430		1.0318		
OCS	n = 6	−29.692	−29.118(1)[b]	0.9807		
	n = 5	−29.865		0.9750		
	n = 4	−26.679		1.0914		
	n = 3	−26.964		1.0799		

[a]Reference 26.
[b]Reference 27.

respectively. Column 5 of Table IV gives the ratios $\chi_{obs.}(^{33}S)/\chi_{calc.}(^{33}S)$ for CS and OCS, while columns 6 and 7 give the values predicted for $\chi_{obs.}(^{33}S)$ of CS_2 by assuming that the ratios that apply for CS and OCS, respectively, also apply to CS_2. Since there is no reason to prefer the results in either column 6 or 7, it seems reasonable to use the mean of the two B3LYP/aug-cc-pV6Z values, that is, to assume $\chi_0(^{33}S)$ = −17.1(6) MHz as the best estimate for the CS_2 free molecule coupling constant.

Correction of the observed ^{33}S nuclear quadrupole coupling constant of the isotopologues $H_3{}^{15}N\cdots{}^{33}S{=}C{=}S$ and $H_3{}^{15}N\cdots S{=}C{=}{}^{33}S$ for the effects of the efg at the appropriate S due to the presence of the NH_3 subunit can be made readily from the observations collected in Table V. Table V lists the two (inner and outer) ^{33}S coupling constants for the isotopologue $H_3N\cdots{}^{33}S{=}C{=}{}^{33}S$ when calculated at the B3LYP/aug-cc-pVnZ (n = 3–6) levels of theory. Also given in Table V is the mean of these two coupling constants (column 4) and the corresponding values of the calculated coupling constant for the free molecule $^{33}S{=}C{=}S$ in column 5. We note that the difference between columns 5 and 6 is negligible in the present context, that is, the mean of the inner and outer ^{33}S coupling constants of $H_3N\cdots{}^{33}S{=}C{=}{}^{33}S$ is essentially identical to the ^{33}S coupling constant of the free CS_2 molecule. Since the calculations refer to the C_{3v} equilibrium geometry of $H_3N\cdots{}^{33}S{=}C{=}{}^{33}S$, in which there is no zero-point oscillation of the subunits ($\theta = \phi = 0$), this analysis suggests that the mean of the values of the $\chi_{obs.}(^{33}S)$ for $H_3{}^{15}N\cdots{}^{33}S{=}C{=}S$ and $H_3{}^{15}N\cdots S{=}C{=}{}^{33}S$ should be used as $\chi_{aa}^{corr}(^{33}S)$ (i.e., corrected for the efg due to NH_3) when deriving $\langle\cos^2\phi\rangle$ from Eq. (4). Thus, $\chi_{aa}^{corr}(^{33}S) = (-16.14 - 16.50)/2 = -16.32$ MHz. When $\chi_{aa}^{corr}(^{33}S)$ = −16.32 MHz and $\chi_0(^{33}S)$ = −17.1(6) MHz are used in Eq. (4), the result is $\phi_{av} = \cos^{-1}\langle\cos^2\phi\rangle^{\frac{1}{2}} = 10(3)°$.

A comparison of $\chi_0(^{14}N)$ for the free ammonia molecule with that of $H_3N\cdots S{=}C{=}S$ when both are calculated at the B3LYP/aug-cc-pVnZ level of theory shows a constant difference of only ~0.04 MHz for n in the range of 3–6. For

TABLE V. ^{33}S nuclear quadrupole coupling constants of $H_3N\cdots{}^{33}S{=}C{=}{}^{33}S$ and $^{33}S{=}C{=}S$ calculated at the B3LYP /aug-cc-pV(n + d)Z level of theory.

Basis	$H_3N\cdots{}^{33}S_i{=}C{=}{}^{33}S_o$ $\chi(^{33}S_i)$ (MHz)[a]	$\chi(^{33}S_o)$ (MHz)[b]	$\chi(^{33}S_{av})$ (MHz)[c]	$^{33}S{=}C{=}S$ $\chi(^{33}S_{CS_2})$ (MHz)	Difference $\{\chi(^{33}S_{av}) - \chi(^{33}S_{CS_2})\}$ (MHz)
n = 3	−14.167	−15.540	−14.854	−14.966	0.112
n = 4	−14.198	−15.539	−14.869	−14.934	0.065
n = 5	−16.230	−17.646	−16.938	−16.958	0.020
n = 6	−16.055	−17.494	−16.775	−16.845	0.070

[a]Subscript i = inner.
[b]Subscript o = outer.
[c]Subscript av indicates the average of i and o values.

example, the two values calculated at the n = 6 level are −4.438 MHz for NH_3 and −4.401 MHz for $H_3N\cdots S{=}C{=}S$. Thus, the free molecule value[19] $\chi_0(^{14}N)$ = −4.090 MHz (included in Table V) can be corrected to $\chi_0^{corr}(^{14}N)$ = −4.090 − 0.04 MHz = −4.13 MHz, which when used in Eq. (4) with the observed value $\chi(^{14}N)$ = −2.846(4) MHz for $H_3{}^{14}N\cdots SCS$ gives $\theta_{av} = \cos^{-1}\langle\cos^2\theta\rangle^{\frac{1}{2}}$ = 27°, with a liberal range of ±3° assumed from consideration of the related average $\langle\theta^2\rangle^{1/2}$ determined in Sec. III D.

D. Intermolecular stretching and bending force constants

There are two measures of the strength of the interaction of NH_3 and S=C=S to form the complex $H_3N\cdots S{=}C{=}S$. The first is the dissociation energy D_e. The second is the intermolecular stretching force constant k_σ, which is a measure of the resistance to an infinitesimal displacement along the weak bond direction. Millen[28] showed that k_σ is simply related to the centrifugal distortion constant D_J and the rotational constant B_0 for symmetric-top complexes such as $H_3N\cdots S{=}C{=}S$ by

$$k_\sigma = \left(16\pi^2\mu B_0^3/D_J\right)\left\{1-(B_0/B^{NH_3})-(B_0/B^{CS_2})\right\}, \quad (5)$$

where μ is as defined earlier and B^{NH_3} and B^{CS_2} are the rotational constants of the separate components. The derivation of Eq. (5) assumes rigid subunits and ignores terms beyond quadratic in the potential energy function. The values of k_σ for the symmetric-top isotopologues calculated using Eq. (5) with the ground-state spectroscopic constants of the complex (Table I) and the separate components (Table II) are included in Table III. The mean is 3.95(5) N m^{-1}, and is slightly smaller than that (k_σ = 5.3 N m^{-1}) recently calculated[13] at the CCSD(T)/aug-cc-pVTZ level of theory. The difference arises in part because, strictly, equilibrium rather than zero-point spectroscopic constants should be used in Eq. (5).

The intermolecular bending force constant $k_{\theta\theta}$ can also be estimated from centrifugal distortion constants. If the model for the motion of the subunits discussed in Sec. III B and illustrated in Fig. 3 is assumed and each subunit is again described as a two-dimensional isotropic harmonic oscillator, the quadratic force constant $k_{\theta\theta}$ is related to the centrifugal distortion constant $D = 2D_J + D_{JK}$ by[29]

$$k_{\theta\theta} = (2h/D)B_0^2\left\{B_0/B^{NH_3}-1\right\}^2. \quad (6)$$

Equation (6) applies to symmetric-top isotopomers of $H_3N\cdots S{=}C{=}S$ only. Values of $k_{\theta\theta}$ estimated by using $D = 2D_J + D_{JK}$ and B_0 from Table I and B^{NH_3} for the free ammonia molecule (Table II) are included in Table III for the symmetric-top isotopologues of the complex. The errors transmitted from those in the centrifugal distortion constants when using Eq. (6) are small, and presumably differences in $k_{\theta\theta}$ outside the quoted errors arise mainly from limitations of the model used to obtain Eq. (6).

An alternative to the angle $\theta_{av} = \cos^{-1}\langle\cos^2\theta\rangle^{\frac{1}{2}}$ defined in Sec. III B is $\langle\theta^2\rangle^{1/2}$, which can be estimated from $k_{\theta\theta}$. It has been shown that $\langle\theta^2\rangle$ is related to $k_{\theta\theta}$ using the following expression:[30]

$$\langle\theta^2\rangle = (h/2\pi)\left(k_{\theta\theta}I_b^{NH_3}\right)^{-1/2}, \quad (7)$$

if the motion of the NH_3 subunit in the symmetric-top complex can be treated as a two-dimensional isotropic oscillator. Using the values of $I_b^{NH_3}$ implied by the monomer rotational constants given in Table II, we estimate that the mean value $\langle\theta^2\rangle^{1/2}$ = 31.0(1)° for the ten symmetric-top isotopologues of $H_3N\cdots S{=}C{=}S$. If $\langle\theta^2\rangle^{1/2}$ is not too large, we can assume that $\langle\theta^2\rangle^{1/2} \approx \cos^{-1}\langle\cos^2\theta\rangle^{\frac{1}{2}}$, which is consistent with the choice of the range for the angle θ_{av} = 27(3)° used in connection with the geometry determination discussed in Sec. III B.

IV. DISCUSSION

The ground-state rotational spectra of 14 isotopologues of the complex $H_3N\cdots S{=}C{=}S$ formed between ammonia and carbon disulphide in the gas phase have been analyzed to yield rotational constants, centrifugal distortion constants, and, where appropriate, ^{14}N and ^{33}S nuclear quadrupole coupling constants. The nature of the spectra and the changes that accompany the isotopic substitutions establish that the complex is a symmetric-top molecule of C_{3v} symmetry with the atoms in the order $H_3N\cdots S{=}C{=}S$. Thus, the intermolecular bond is between the non-bonding electron pair on N and the electrophilic region near S on the C_3 axis of the complex, that is, it is the type of non-covalent interaction now referred to as a chalcogen bond. A simple model used to account for the contribution of the angular oscillations to the zero-point motion leads to the distance $r(N\cdots S)$ = 3.338(10) Å, which is, as expected, longer than the equilibrium value 3.304 Å obtained from a recent *ab initio* calculation[13] at the CCSD(T)/aug-cc-pVTZ level. Interpretations of the centrifugal distortion constants of the complex lead to the intermolecular stretching and bending force constants k_σ = 3.95(5) N m^{-1} and $k_{\theta\theta} = 4.63(5)\times10^{-21}$ J, which indicate that the intermolecular chalcogen bond is weak. For comparison, values of k_σ = 12.2 and 12.7 N m^{-1} and $k_{\theta\theta} = 27.5\times10^{-21}$ and 42.2×10^{-21} J have been reported by these methods[30] for the hydrogen-bonded complex $H_3N\cdots HCN$ and the halogen-bonded complex $H_3N\cdots Cl_2$, respectively. The value of $\langle\theta^2\rangle^{1/2}$ = 31° determined from $k_{\theta\theta}$ has been shown to be consistent with the angular oscillation $\theta_{av} = \cos^{-1}\langle\cos^2\theta\rangle^{\frac{1}{2}}$ = 27(3)° of the NH_3 subunit deduced from the ^{14}N nuclear quadrupole coupling constant.

SUPPLEMENTARY MATERIAL

See supplementary material for outputs from the final cycle of the PGOPHER fits of transitions of each isotopologue investigated and from the GAUSSIAN calculations.

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